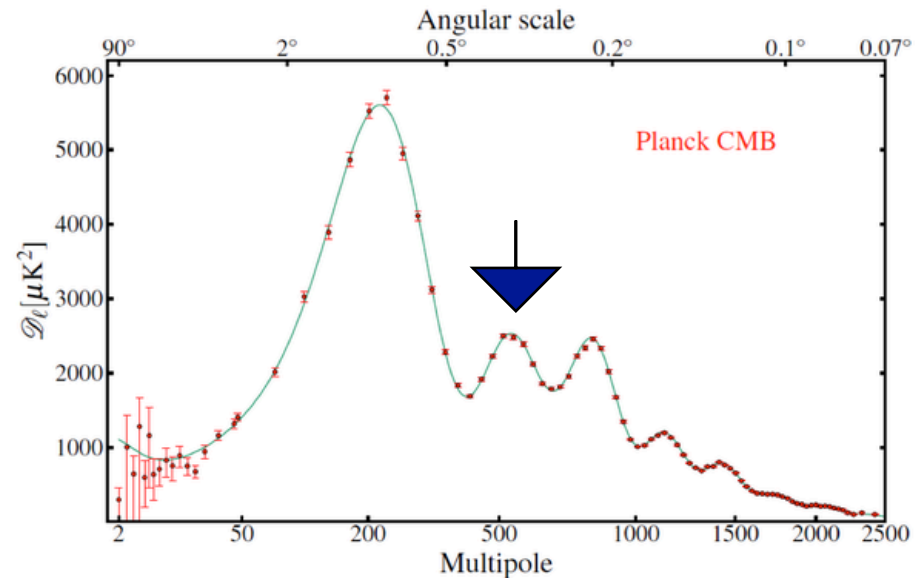
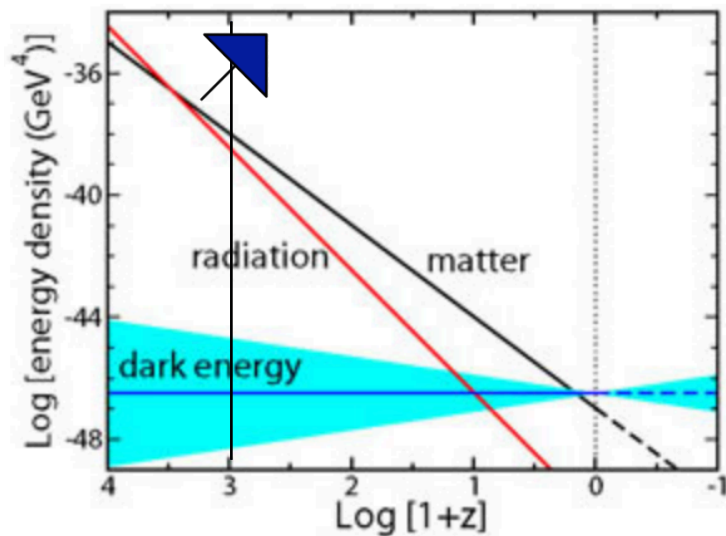


Galaxies and the nature of dark matter & gravity

Benoit Famaey

CNRS - Observatoire astronomique de Strasbourg

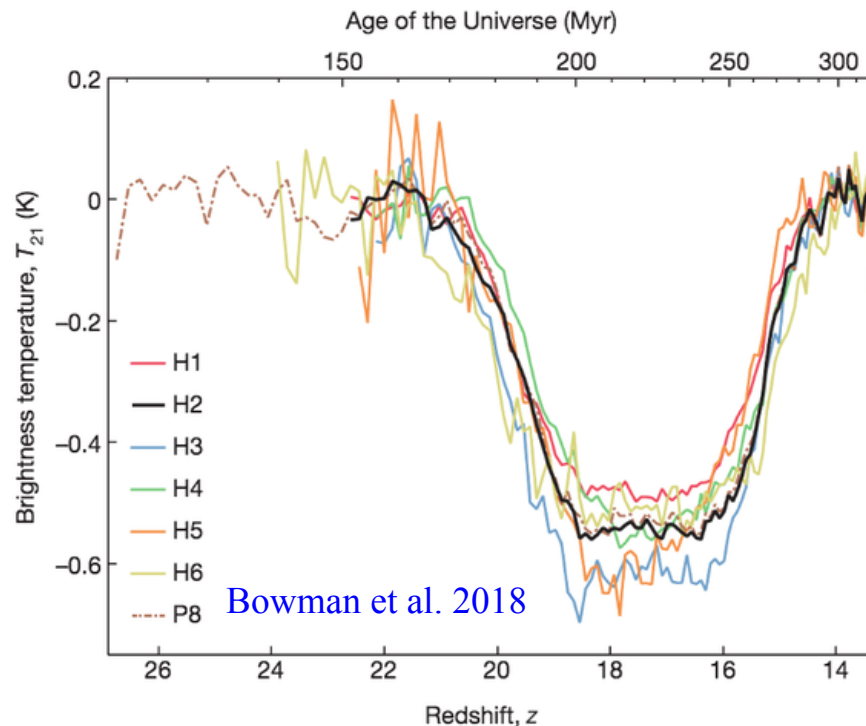
The need for cosmic Dark Matter



- CMB + other large scale probes \Rightarrow concordance Λ CDM model
- DM = **collisionless** and dissipationless fluid of stable elementary particles which interact with each other and with baryons (almost) entirely through gravity, & non-relativistic (**cold**) at matter-radiation equality to form structures down to small scales

Cosmological tensions and the nature of DM

- The **Hubble tension**? No one is really sure what is going on (e.g., Di Valentino et al. 2021)
- The **EDGES anomaly**: no one knows either, potentially a fluke? If not, might have consequences on the nature of DM



- Cosmic dawn absorption feature at $z \sim 17$
- **Factor of 2 too large**
=> fluke?
or background temp. higher at these wavelengths ?
or gas cooler ?



« Small-scale » tensions and the nature of DM

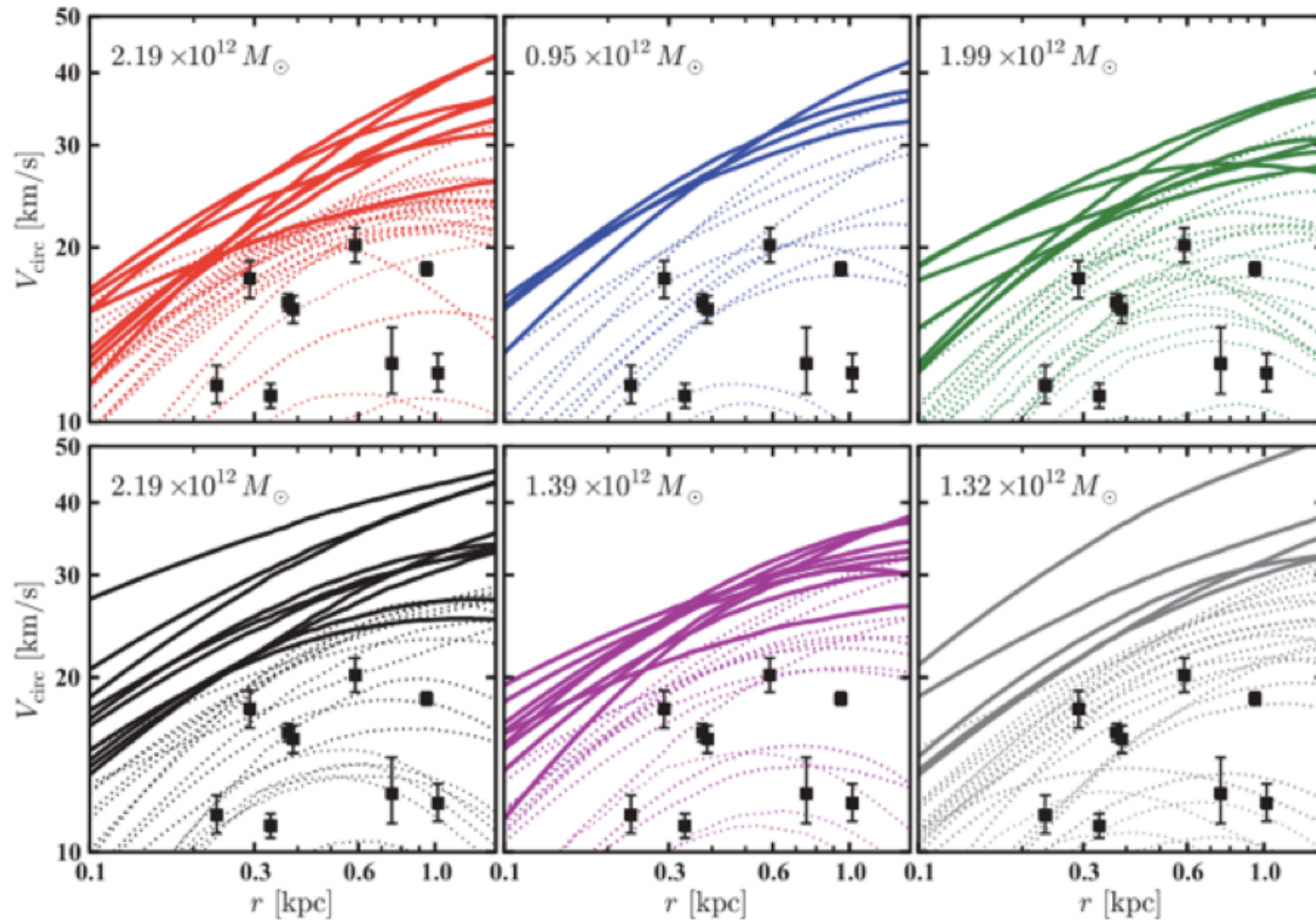
- Galaxies in non-linear ($|\delta| \gg 1$) regime of structure formation
- It is **hard** because of the importance of baryonic physics (feedback!)
- Simulations have made **huge improvements** at forming more realistic galaxies, but some tensions persist...
- Could the problem be **fundamental**, i.e. mostly the nature of DM in the model?
- Typically two types of cosmological galaxy formation sims:
 - **Large box**: EAGLE, IllustrisTNG, HorizonAGN, ...
 - **Zoom-in**: APOSTLE, NIHAO, FIRE-2, Auriga, ...
(can also make constrained simulations like HESTIA for LG)



« Small-scale » tensions and the nature of DM

- Too-big-to-fail (TBTf)
- Tightness of baryonic Tully-Fisher relation (BTFR)
- Failed feedback problem (FFP)
- Diversity of rotation curves problem (modern core-cusp)
- Hot orbits problem
- Fast bar problem
- Satellites phase-space correlation problem (planes of satellites)

Dwarf spheroidal galaxies: Too-big-to-fail



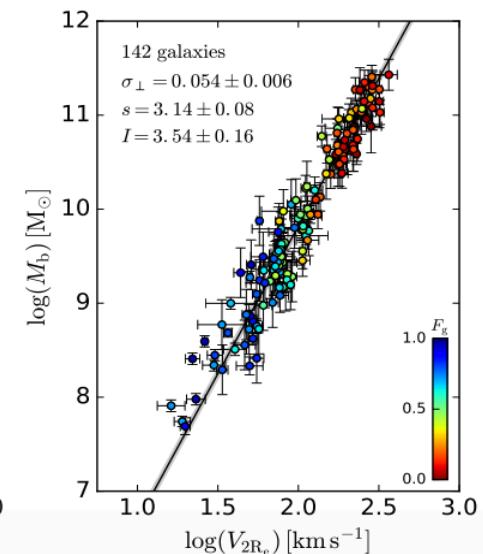
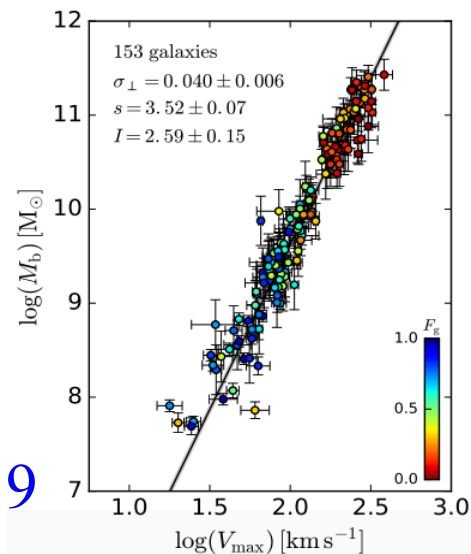
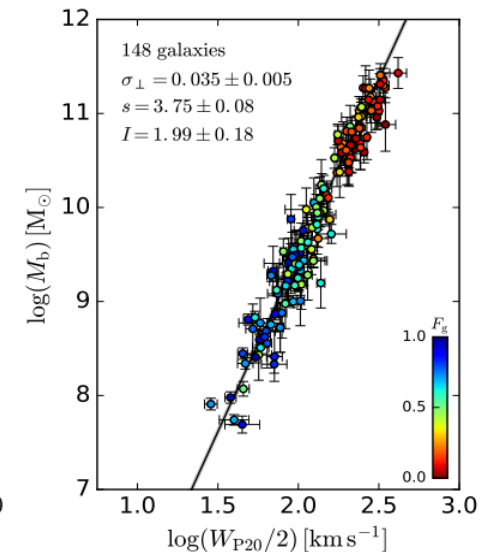
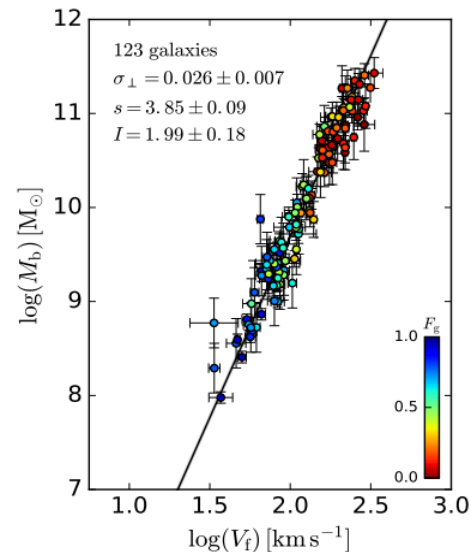
Boylan-Kolchin et al. 2012

Disk galaxies: the baryonic Tully-Fisher relation (BTFR) and its scatter

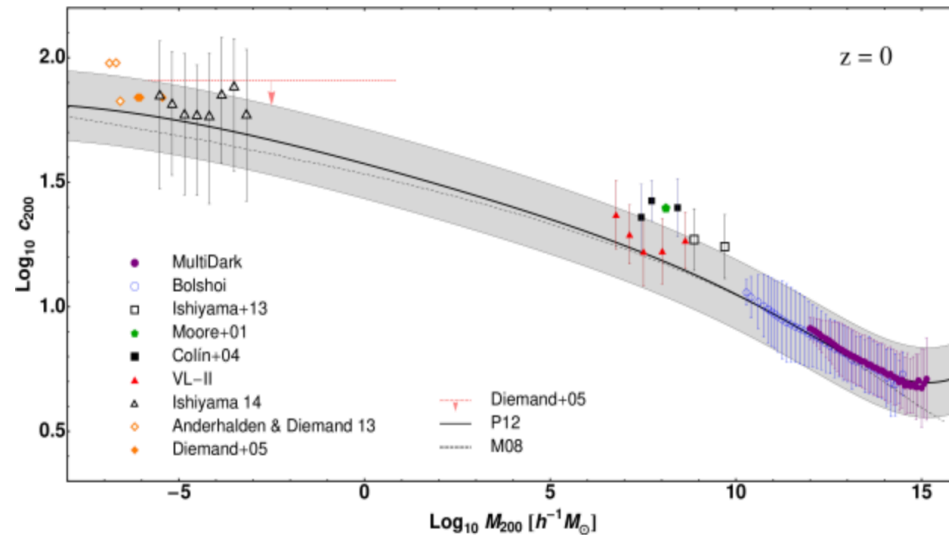
- $\log M_b = \alpha \log V_f - \log \beta$
- $\alpha \approx 3.9 - 4$

- Intrinsic scatter
 ~ 0.025 dex

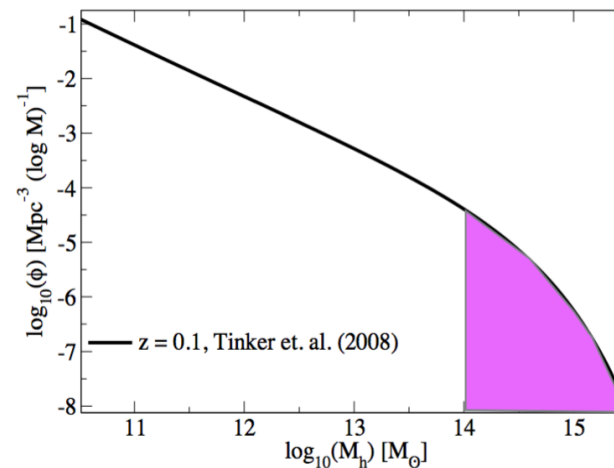
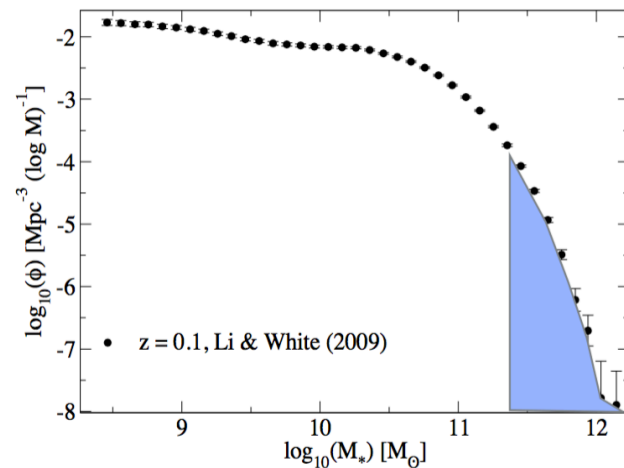
Lelli et al. 2019



Halo scaling relations and abundance matching

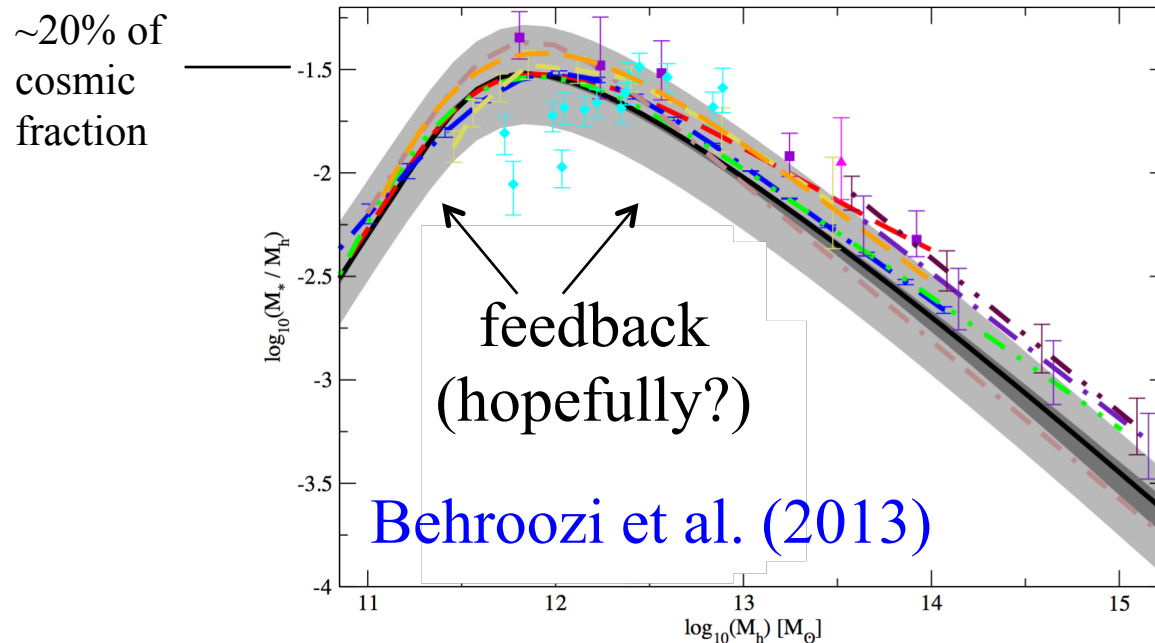


Halo mass-concentration relation
(with some scatter of ~ 0.1 dex)



Match stellar mass
function to halo
mass function by
assigning $n(>M^*)$
to $n(>M_h)$

Stellar-to-halo mass relation (SHMR)



Typical scatter ~ 0.15 dex

⇒ Adding the gas, the intrinsic BTFR scatter **cannot go below 0.05 dex**

Twice too high!

The scatter, residual correlations and curvature of the SPARC baryonic Tully–Fisher relation

Harry Desmond^{1,2*} (2017)

¹Kavli Institute for Particle Astrophysics and Cosmology, Physics Department, Stanford University, Stanford, CA 94305, USA

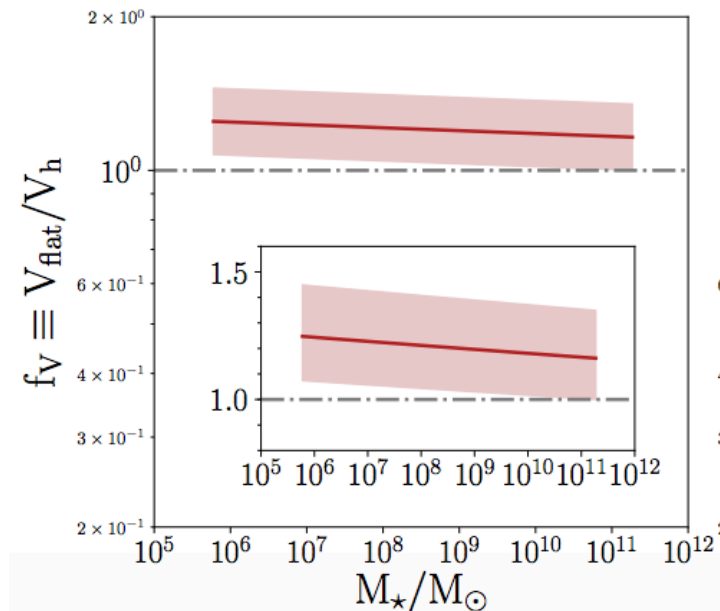
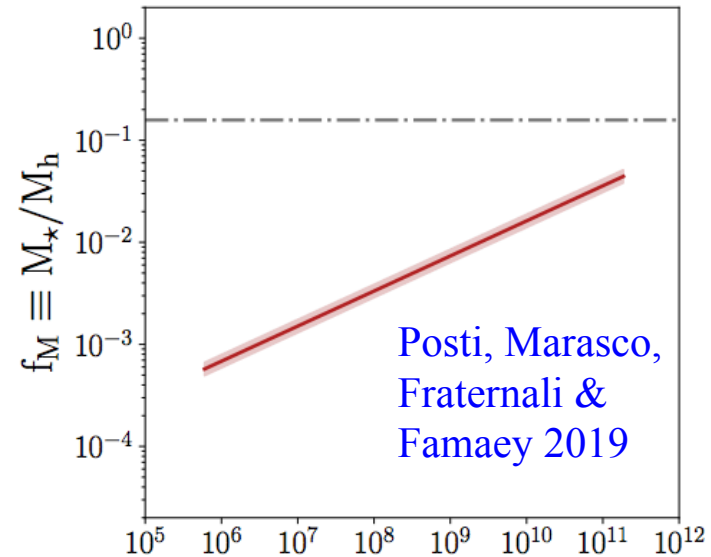
calculate the statistical significance of these results in the framework of halo abundance matching, which imposes a canonical galaxy–halo connection. Taking full account of sample variance among SPARC-like realisations of the parent halo population, we find the scatter in the predicted BTFR to be **3.6 σ too high.**

The failed feedback problem

Reverting the problem: constraining simultaneously M^*/M_{vir} & $V_{\text{flat}}/V_{\text{vir}}$ to fit the high-mass end of Tully-Fisher (together with M^* -size and M^* -j)

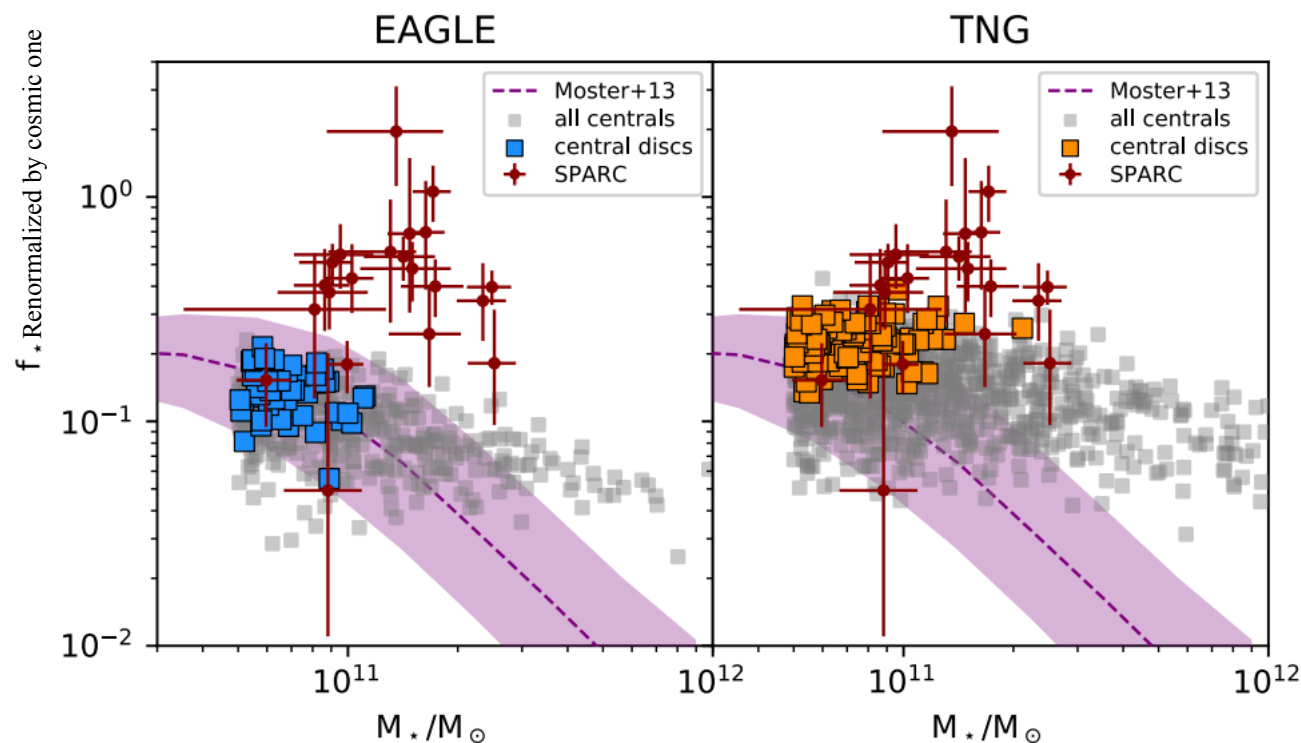
$\Rightarrow M^*/M_{\text{vir}}$ **grows linearly with mass** for disk galaxies, contrary to abundance matching expectations

\Rightarrow **failed feedback** in massive spirals



The failed feedback problem

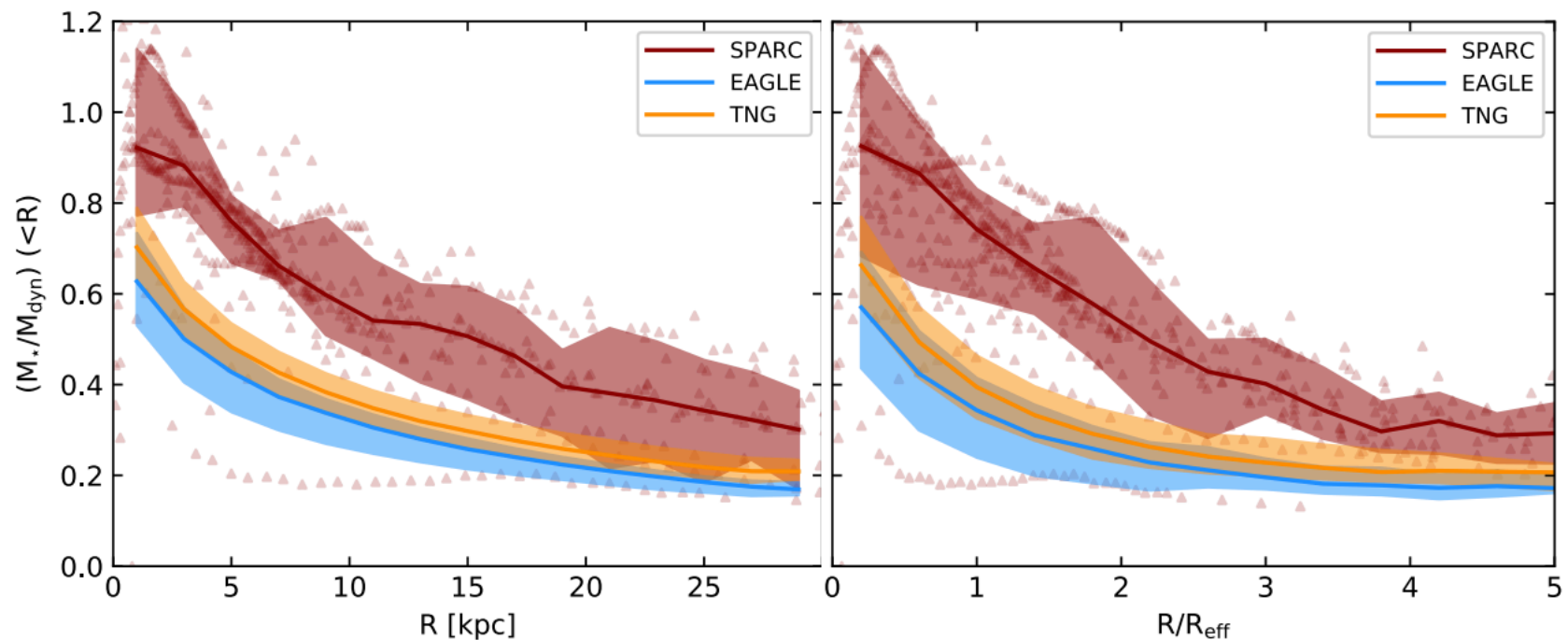
Particle DM mass resolution $< 10^7 M_{\text{sun}}$, **EAGLE** and **Illustris TNG100** allow for a fair evaluation of the behavior of massive disks in simulations



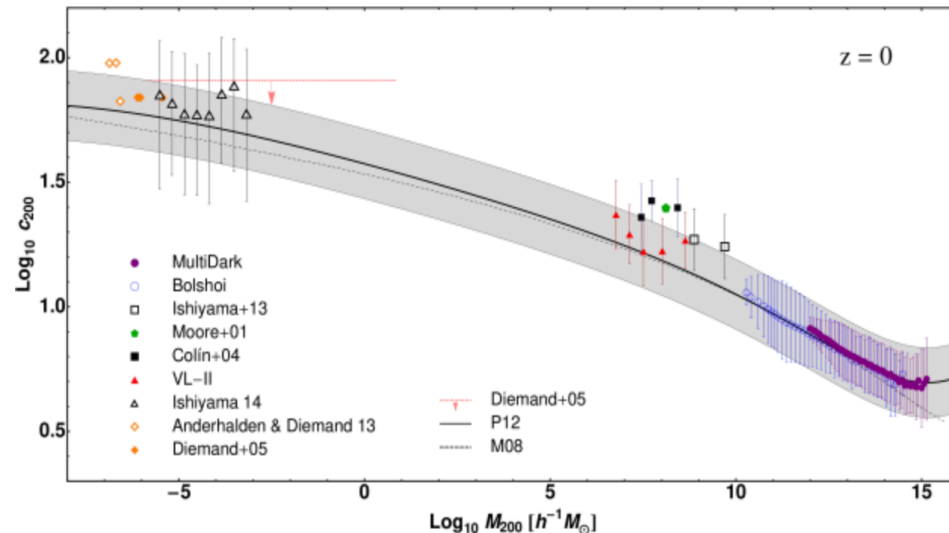
Marasco, Posti, Oman, Famaey, Cresci & Fraternali 2020

The failed feedback problem

Simulated halos hosting massive disks are too inefficient at converting their baryons into stars, through **too efficient feedback**, AND they have undergone halo contraction because of apparently **not efficient enough feedback**...



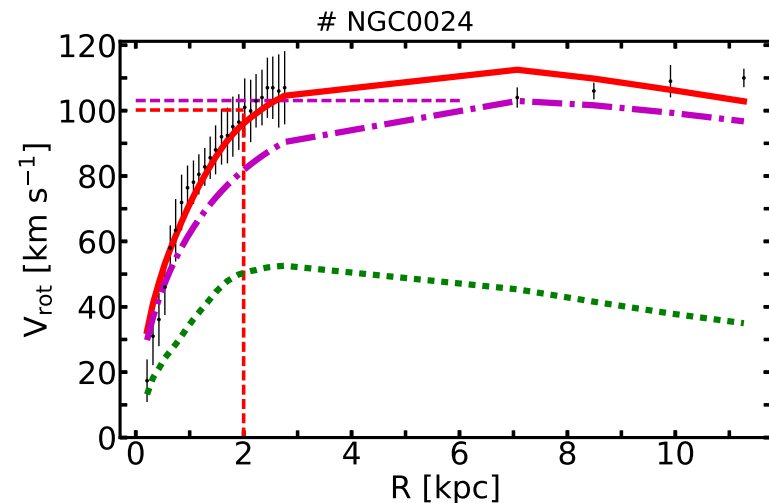
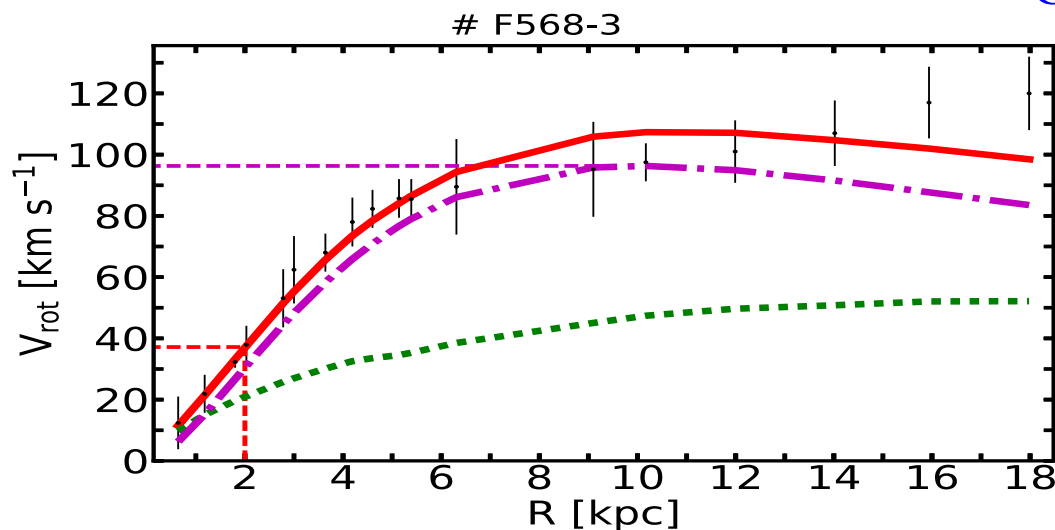
The BTFR twin paradox



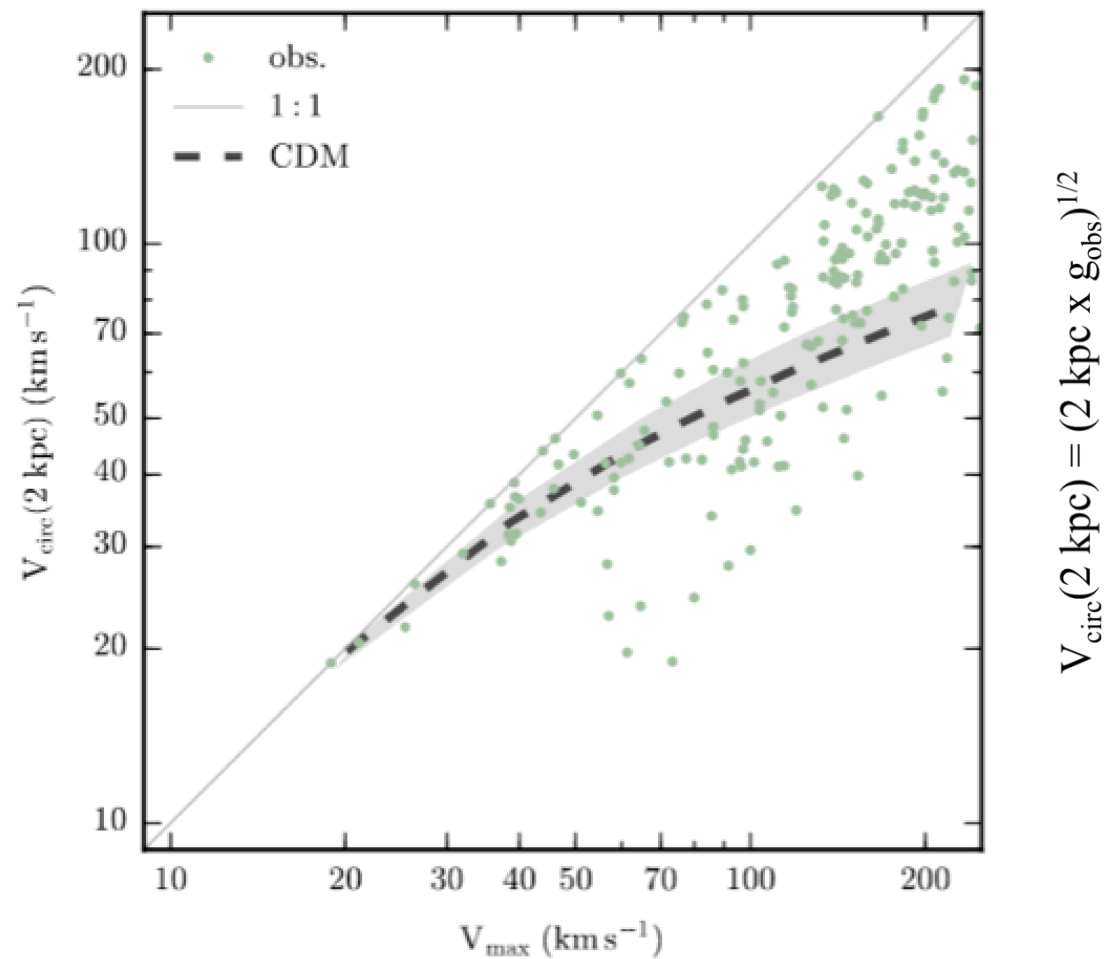
Dark matter halos are (almost) a one-parameter family (driven by mass)

\Rightarrow At the same V_{flat} , why so different profiles??

Ghari, Famaey, Laporte & Haghi 2019

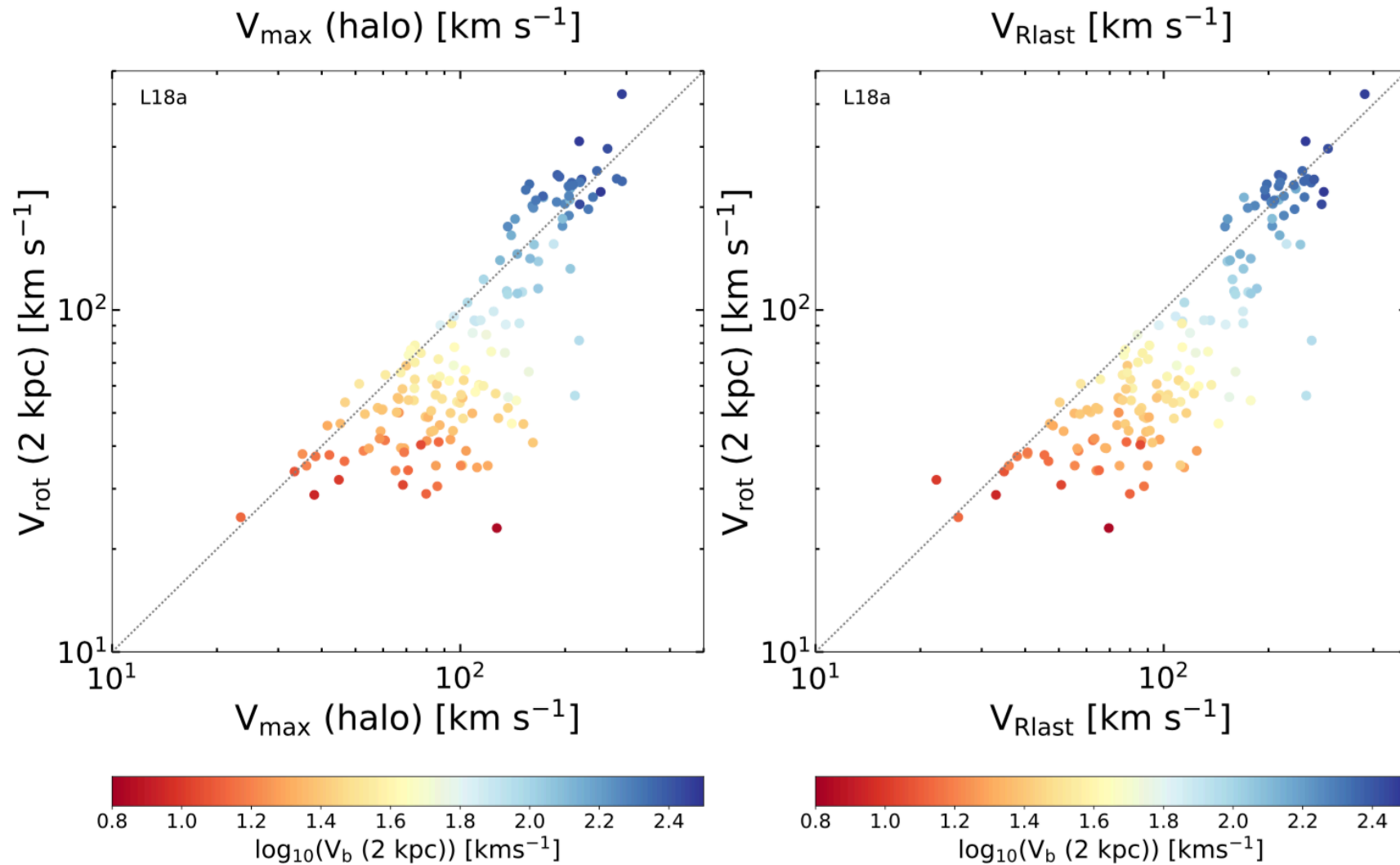


The diversity problem



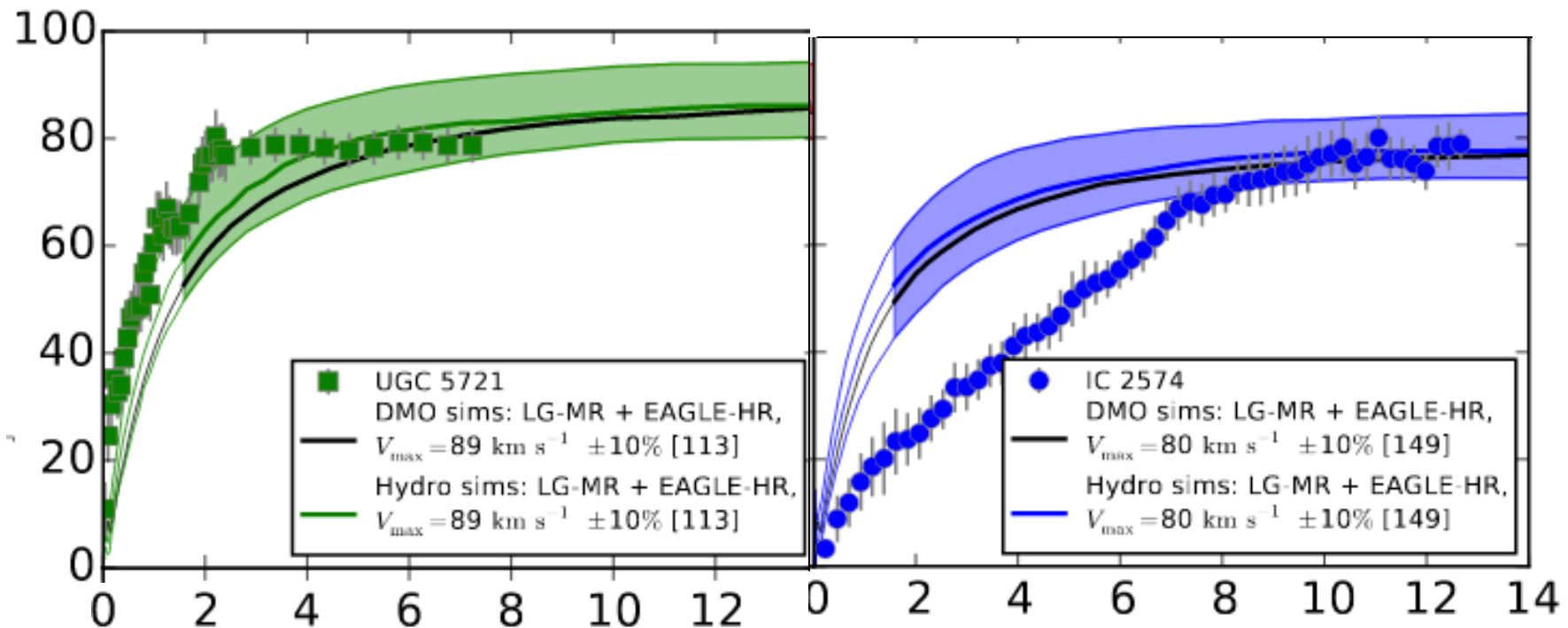
Oman et al. 2015, Bullock & Boylan-Kolchin 2017

Diversity driven by the baryons



Ghari, Famaey, Laporte & Haghi (2019)

The diversity problem or the modern core-cusp problem

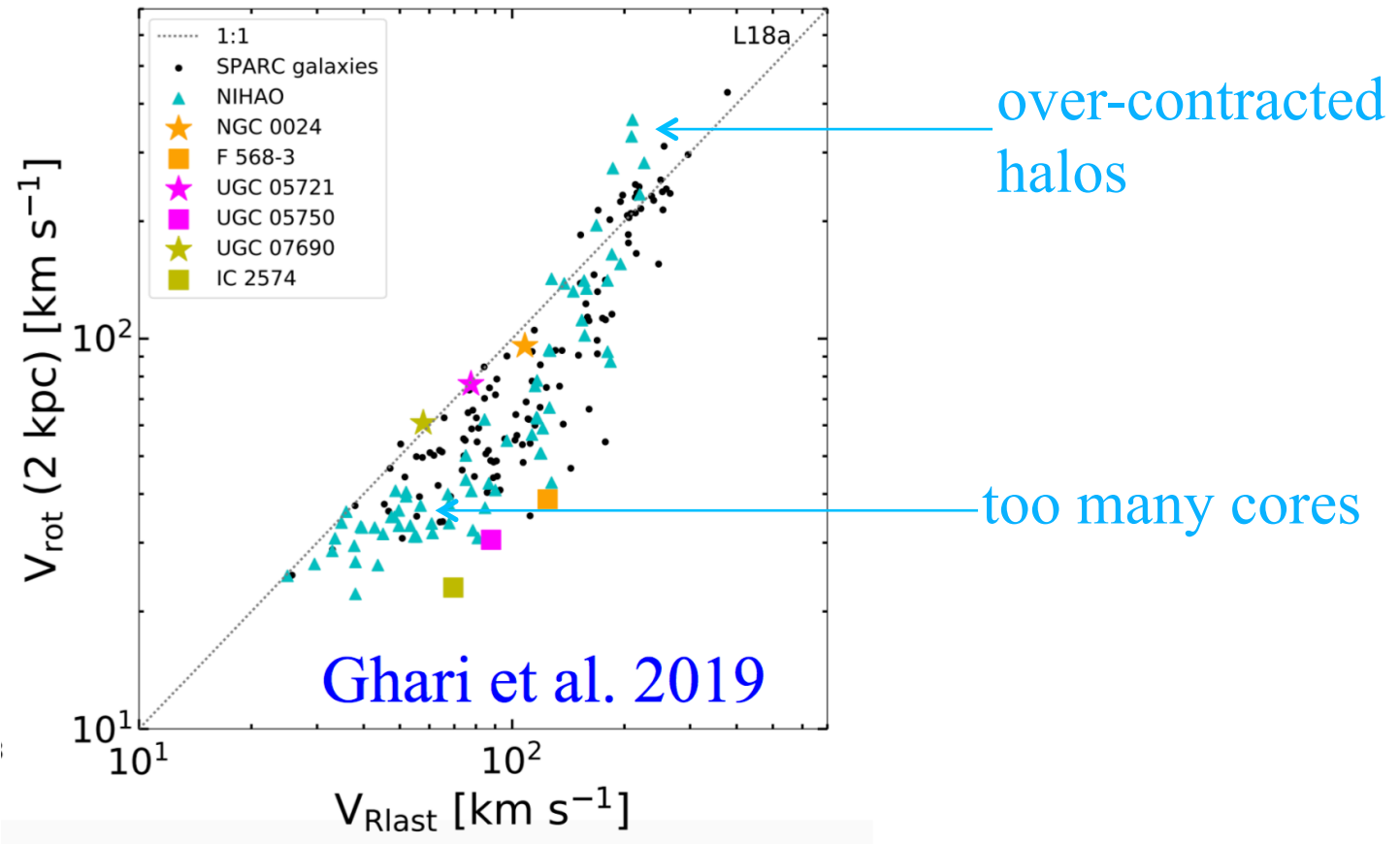


APOSTLE/EAGLE simulations
=> cannot form cores

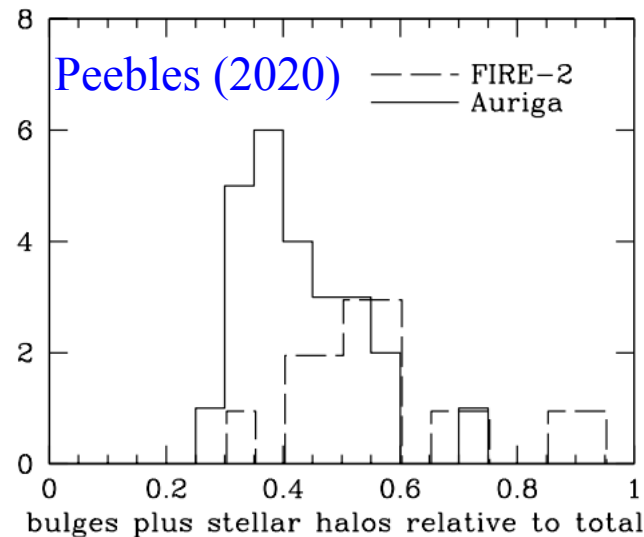
Oman et al. 2015

The diversity problem or the modern core-cusp problem

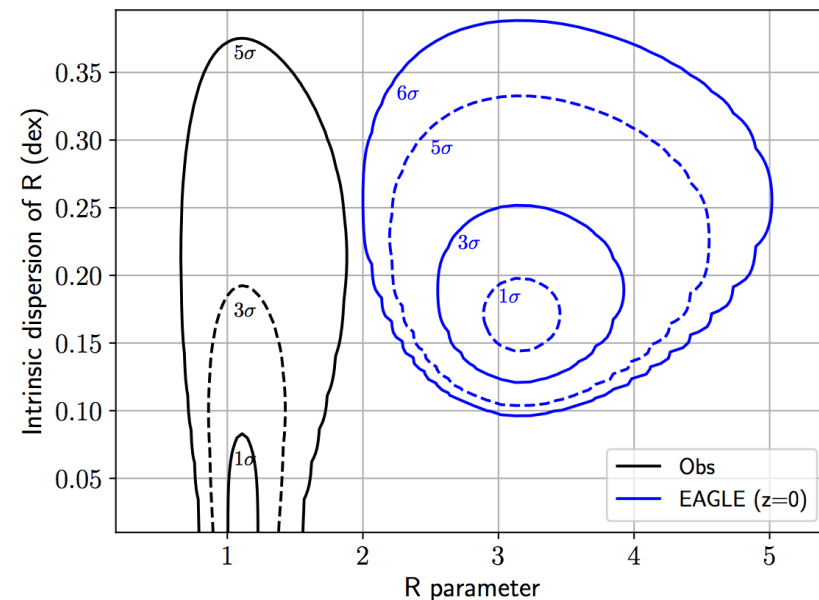
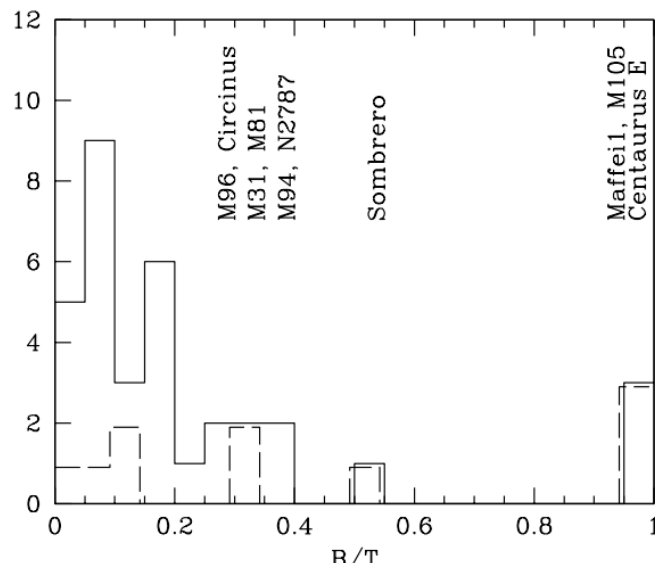
NIHAO has a rather extreme feedback recipe,
leading to **too many cores** at low masses :



The hot orbits problem and the fast bar problem

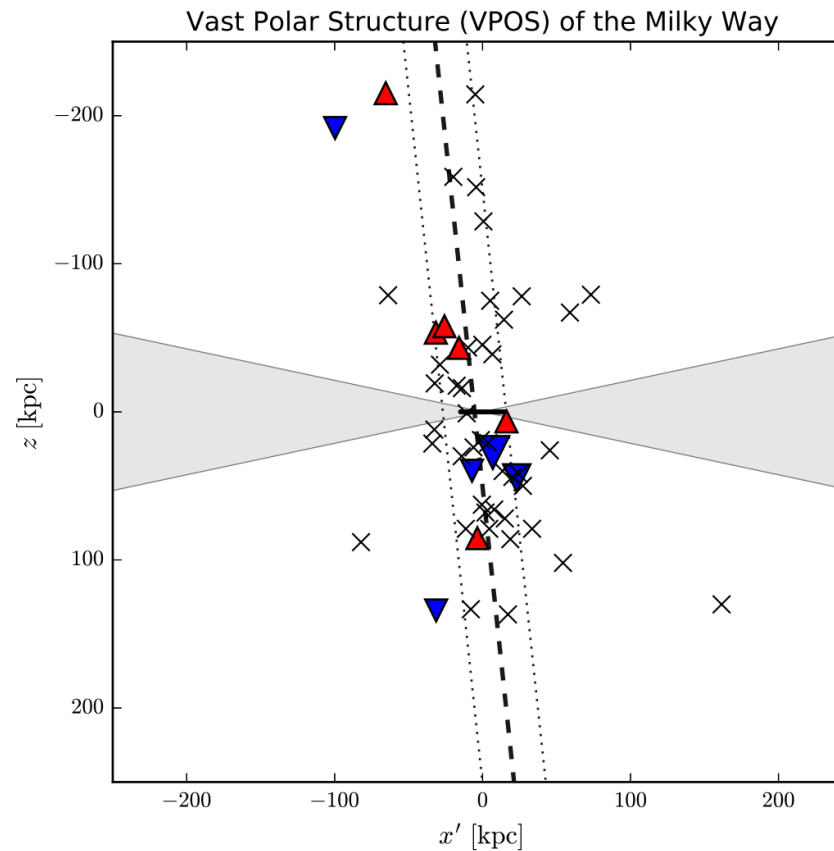


- Most local disk galaxies are nearly **bulgeless** with light stellar halos
- **70% are barred** at $M_* \sim 10^9\text{-}10^{10}M_{\text{sun}}$ (Erwin 2018)
- Bars are **fast** $R_{\text{CR}}/R_{\text{bar}} < 1.4$ (Aguerri et al. 2015)



Roshan et al. (2021)

The satellites phase-space correlation problem



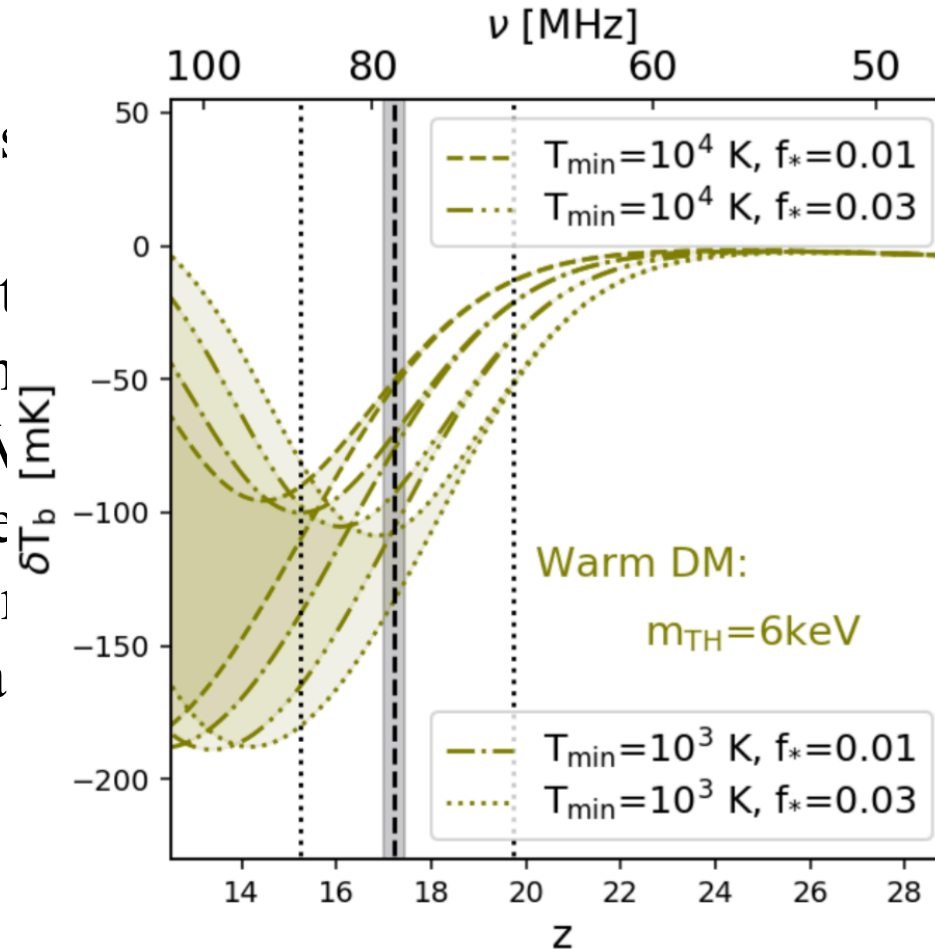
Pawlowski (2018)

Warm dark matter?

The simplest
CDM often as

What about st

- Dam
- (1 keV
- lowe
- To c
- forma



have to be cold?
o a few TeV

DM of a few keV?
-streaming scale

helps TBTF
which prevents the
esnt help diversity

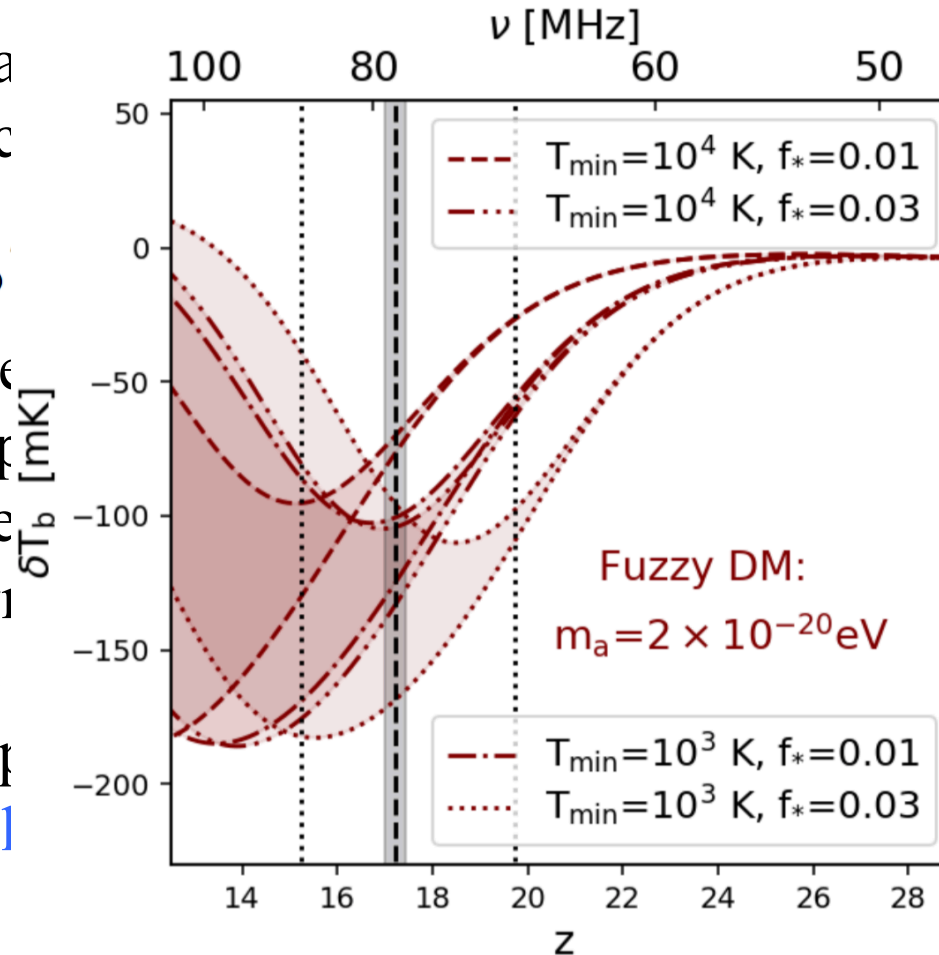
Schneider (2018): delayed formation of small-scale halos in contradiction with EDGES timing for $m < 7 \text{ keV}$ (but at higher masses, **cannot solve any small-scale tension !**)

Fuzzy dark matter?

An idea that ga
DM might be c

$$\lambda = h/(m_b$$

- Above
- **Dam**
- Create
- of magi
- These
- **orbits**
- Not cl



Witten (2017) that
oglie wavelength:
 $100 \text{ km s}^{-1} / v$

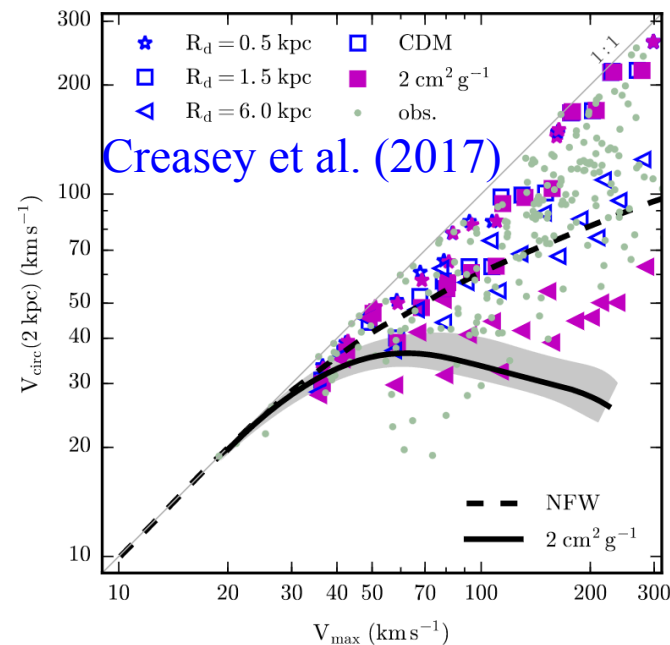
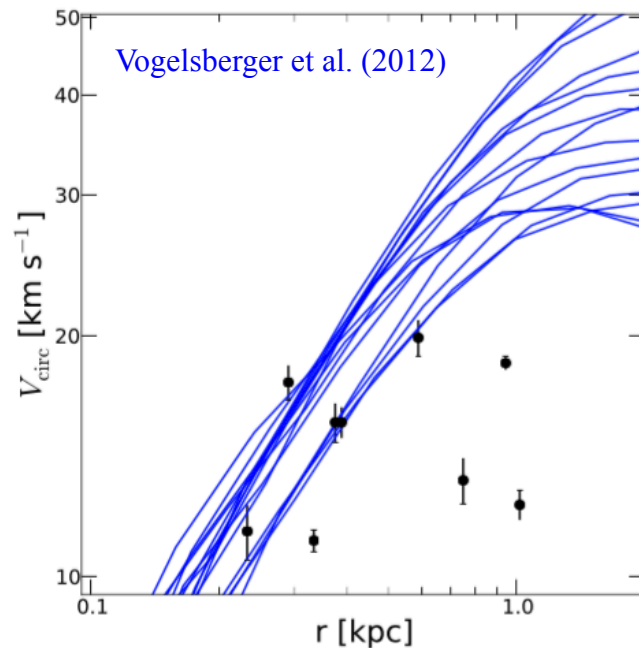
it is **different**
' 10^{-22} eV) $^{-4/3} M_{\odot}$
iction by one order
ale fluctuations)
roblem, maybe **hot**
tness
sity problem

Schneider (2018): delayed formation of small-scale halos in
contradiction with EDGES timing for $m < 10^{-20} \text{ eV}$ (but at higher
masses, **cannot solve any small-scale tension !**)

Self-interacting dark matter?

The 2nd simplest modif. of DM: **does it really have to be collisionless?**
Self-interactions have little effect on the matter power spectrum, but can drastically change the DM profiles in relaxed clusters!

Self-interacting cross-sections $\sigma/m = 1-10 \text{ cm}^2/\text{g}$ can have a drastic effect on halo profiles => can solve **TBTF, diversity, and fast bar problem**



Nothing on hot orbits, and might make **FFP worse!** (Sameie et al. 2021)

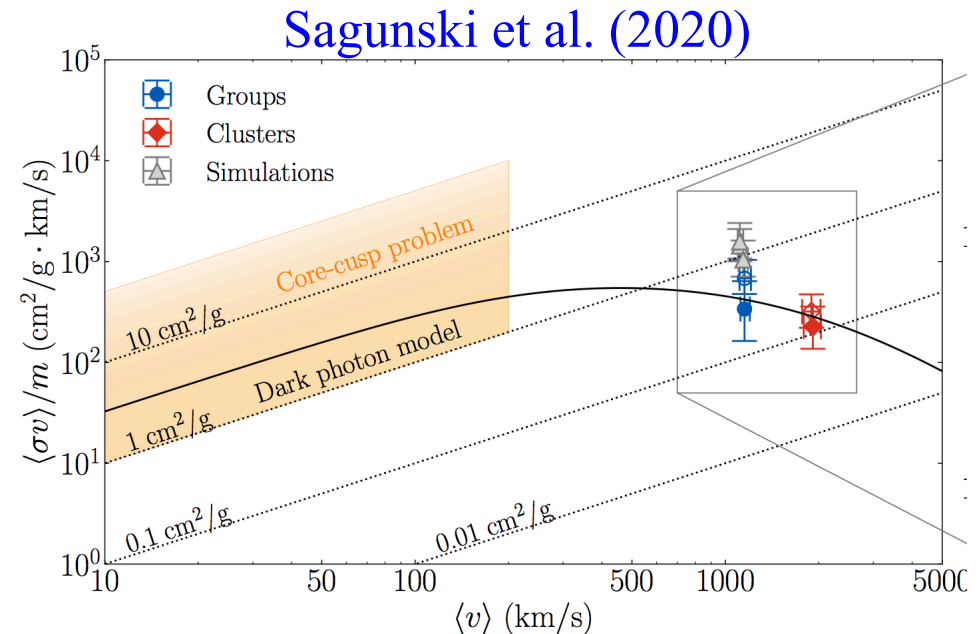
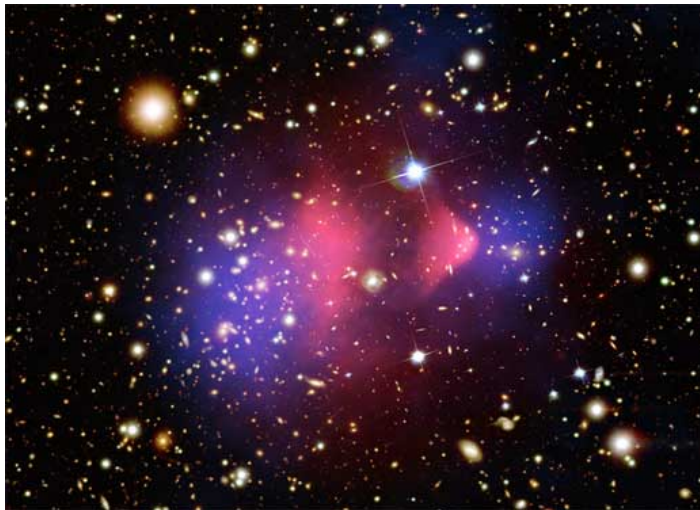
Self-interacting dark matter?

Conflicting constraints with galaxies coming from galaxy clusters:

Colliding clusters (bullet) $\Rightarrow \sigma/m < 0.7 \text{ cm}^2/\text{g}$ (Randall et al. 2008)

Strong lensing of cluster centers $\Rightarrow \sigma/m < 0.065 \text{ cm}^2/\text{g}$ (Andrade et al. 2021)

Cannot solve any tension on galaxy scales with such cross-sections
 \Rightarrow **velocity-dependent cross-section needed**



Modifying gravity?

$$\begin{aligned} g &= g_N \\ g &= (g_N a_0)^{1/2} \end{aligned}$$

$$\begin{aligned} &\text{if } g \gg a_0 \\ &\text{if } g \ll a_0 \end{aligned}$$

MOND

Milgrom 1983

A characteristic **acceleration scale** present in the BTFR and diversity

$$\nabla \cdot [\mu(|\nabla\Phi|/a_0) \nabla\Phi] = 4\pi G \rho_{\text{bar}} \quad \text{AQUAL: Bekenstein \& M (1984)}$$

or

$$\nabla^2 \Phi = \nabla \cdot [\nu(|\nabla\Phi_N|/a_0) \nabla\Phi_N] \quad \text{QUMOND: Milgrom (2010)}$$

⇒ Getting a **tight and straight BTFR**, solving the **failed feedback** problem and the **diversity** for free

+ **no dynamical friction with DM** implies for instance **faster bars as observed**, and **reduces formation of bulges (hot orbits problem)**



Modifying gravity?

$$g = g_N$$

$$g = (g_N a_0)^{1/2}$$

$$\text{if } g \gg a_0$$

$$\text{if } g \ll a_0$$

MOND

Milgrom 1983

⇒ Question: how to get the right CMB peaks?

Needs a new degree of freedom which decouples from the baryon-photon plasma in time-dependent configurations, and which gives MOND in equilibrium quasi-static configurations

⇒ **RelMOND (Skordis & Zlosnik 2020)**

Modifying gravity?

$$\begin{aligned} g &= g_N && \text{if } g \gg a_0 \\ g &= (g_N a_0)^{1/2} && \text{if } g \ll a_0 \end{aligned}$$

MOND
Milgrom 1983

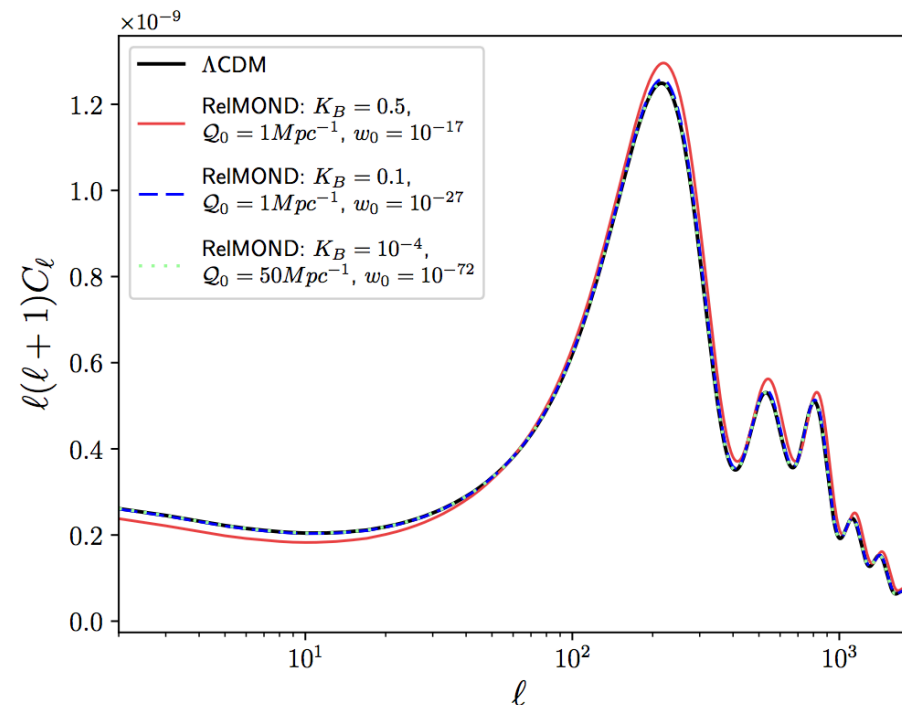


FIG. 1. The CMB temperature angular power spectrum C_ℓ for Λ CDM (black) and relativistic MOND for a number of parameter values (keeping all parameters common to Λ CDM the same). The green (dotted) curve is indistinguishable from Λ CDM.

Modifying gravity?

$$g = g_N$$

$$g = (g_N a_0)^{1/2}$$

$$\text{if } g \gg a_0$$

$$\text{if } g \ll a_0$$

MOND

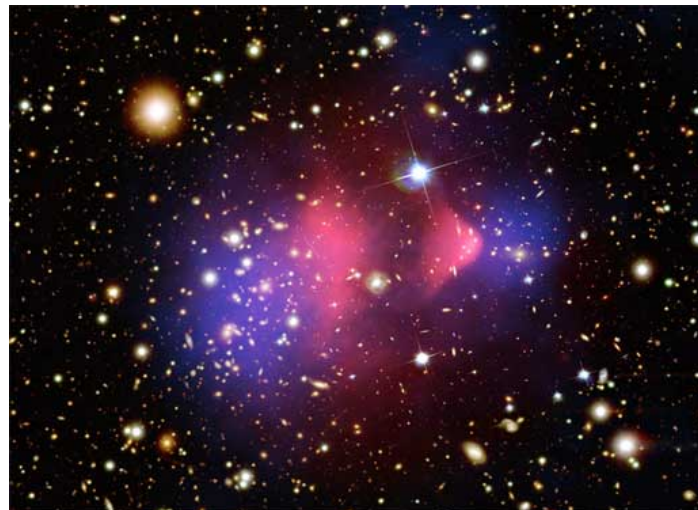
Milgrom 1983

⇒ **Real challenge: non-linear regime and galaxy clusters!**

Intermediate regime of barely virialized systems??

Ultra-diffuse galaxies in clusters immune to the EFE?

(Freundlich, Famaey, et al. in prep.)



??

Baryon-interacting dark matter?

Change from CBE to BTE with two fluids through some long-range interaction (Famaey et al. 2018, 2020)

⇒ second order moments then give a **heat equation** which can resemble the MOND equation if roughly assuming $T \propto \Phi$

$$\frac{3}{2} \left(\frac{\partial}{\partial t} + \vec{u} \cdot \vec{\nabla} \right) \frac{T}{m} + \frac{1}{\rho} P^{ij} \partial_i u_j + \frac{1}{\rho} \vec{\nabla} \cdot \vec{q} = \frac{\dot{\mathcal{E}}}{m}$$

Spherical symmetry+isotropy+no spin+equilibrium (no t dependence) for halo:

$$\vec{\nabla} \cdot \left(m \kappa \vec{\nabla} v^2 \right) = -\rho \frac{\dot{\mathcal{E}}}{m}$$

Two things to fix: **thermal conductivity** and **heating rate**



Thermal conductivity :

$$\kappa = \frac{3}{2} \frac{\rho v^2 t_{\text{relax}}}{m} \quad \text{through some sort of DM self-interactions}$$

Needs a relatively short relaxation time, let's take: $t_{\text{relax}} = \frac{\mathcal{N}}{\sqrt{G\rho}}$

Heating rate :

We want a_0 in the denominator on the l.h.s., hence should be prop. to a_0 , simplest is to take $a_0 v$, and dimensionless dependence on ρ and ρ_b

$$\frac{\dot{\mathcal{E}}}{m} = C a_0 v \frac{\rho_b}{\rho}$$

\Rightarrow We showed it gives rise to right diversity
 \Rightarrow little interaction for CMB, just the right energy exchange for **EDGES**... (simply by putting a_0 scale in the heating rate)



Conclusions on « small-scale » tensions and the nature of DM

- WDM: good for TBTF, not so much for the other challenges, **above ~ 10 keV, does not really solve any challenge**. Perhaps hot orbits if coupled with non-gaussianities
- FDM: good for TBTF and reducing dynamical friction, not so much other challenges such as diversity of RC, **above $\sim 10^{-20}$ eV, does not really solve any challenge**
- SIDM: very promising for diversity! **could make failed feedback at the high mass end worse**, velocity-dependence tightly constrained by galaxy clusters
- MOND: solves quite a few challenges at galaxy scales! But also creates new ones (convoluted relativistic theory, **missing mass in clusters, UDGs in clusters,...**)
- BIDM: not explored very much yet...